

# Factorization of Lepton Radiation in SIDIS

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# Proton Spin Question

- Total spin decomposed as in (Jaffe and Manohar, 1990; Ji, 1997)

$$J = \frac{1}{2} \Delta\Sigma + \Delta G + L_q + L_g$$

- Sum of quark spins do not match proton spin

$$\Delta\Sigma = \Delta u + \Delta d + \Delta s \approx 0.3$$

- Polarized DIS experiments complementary to RHIC to determine  $\Delta G$  (currently  $\approx 0.4$ )
- $L_{q/g}$  related to transverse motion  $\Rightarrow$  TMDs



# JLab, EIC, and SIDIS

- Current Jefferson Lab (JLab) 12 GeV and upcoming Electron Ion Collider (EIC) will perform electron-proton scattering to probe structure of proton
- Structure probed in semi-inclusive deep inelastic scattering (SIDIS) is related to parton distribution functions (PDFs) and transverse momentum dependent PDFs (TMDs)
- Factorization connects measured hadron to PDFs and TMDs
- Kinematics defined in photon-hadron frame under the one-photon approximation

$$d\sigma \propto \left| \begin{array}{c} \text{Diagram} \end{array} \right|^2 \propto L^{\mu\nu}(q, \ell) W_{\mu\nu}(q, p, \dots)$$

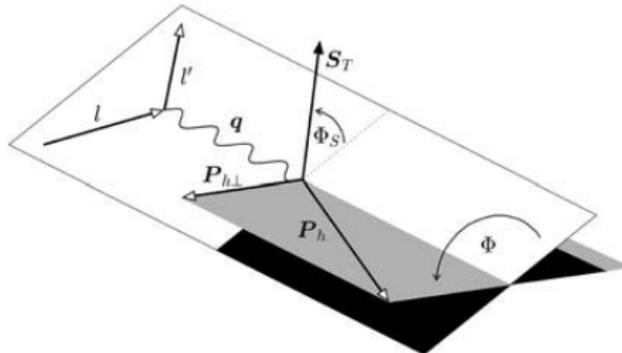
- Collinear factorization separates hadronic portion into two parts: PDFs and fragmentation functions (FFs)

$$E_h E_{\ell'} \frac{d^3\sigma^{\ell A \rightarrow \ell' h X}}{d^3\vec{P}_h d^3\vec{\ell}'} = \frac{1}{2S} \sum_{i,j} \int \frac{dx}{x} \frac{dz}{z^2} \tilde{f}^{i/A}(x) \tilde{D}^{h/j}(z) \hat{\sigma}^{ei \rightarrow e'j}(x, z) \quad (1)$$



# Photon-Hadron Frame

- Standard definitions uses Trento Convention (Bacchetta et al., 2004)



- Radiation of photon from initial state changes the plane modulation
- Unobserved lepton radiation introduces ambiguity, as the radiation changes the unobserved photon's momentum (Liu et al., 2021)
- Unified factorization method to address these issues

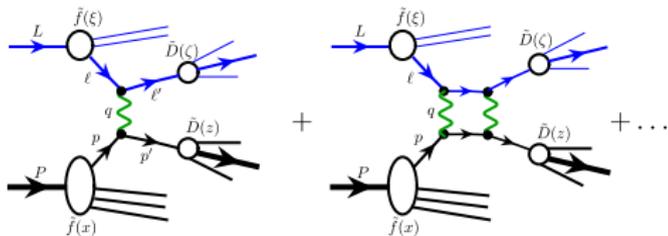


## New Factorization Approach

- Unified factorization scheme (QED and QCD on equal footing) (Liu et al., 2021)
- Leading collinear radiative corrections resummed into lepton distribution functions and lepton fragmentation functions
- Focus on collinear regime at first

$$E_h E_{L'} \frac{d^3 \sigma^{LA \rightarrow L' h X}}{d^3 \vec{P}_h d^3 \vec{L}'} = \frac{1}{2S} \sum_{i,j,n,m} \int \frac{d\zeta}{\zeta^2} \frac{d\xi}{\xi} \frac{dz}{z^2} \frac{dx}{x} \tilde{D}_{m/e}(\zeta) \tilde{f}_{e/i}(\xi) \\ \times \tilde{D}_{h/n}(z) \tilde{f}_{j/h}(x) \hat{H}^{ij \rightarrow nm}(\xi, \zeta, x, z) \quad (2)$$

- Note no structure functions in this expression (which is valid beyond the one-photon exchange)
- Can still extract PDFs and TMDs from scattering and can naturally take care of higher order QED contributions



## Lepton-Hadron (Lab) Frame

- Since the direction of the virtual photon is unknown, two common kinematic frames chosen
- Defining the external momenta in the lab frame using the rapidity and transverse momenta,

$$L^\mu = \left( \sqrt{\frac{S}{2}}, 0^-, \vec{0}_\perp \right), \quad L'^\mu = \left( \frac{L'_T}{\sqrt{2}} e^{y_L}, \frac{L'_T}{\sqrt{2}} e^{-y_L}, \vec{L}'_T \right)$$

$$P^\mu = \left( 0^+, \sqrt{\frac{S}{2}}, \vec{0}_\perp \right), \quad P'^\mu = \left( \frac{P'_T}{\sqrt{2}} e^{y_P}, \frac{P'_T}{\sqrt{2}} e^{-y_P}, \vec{P}'_T \right)$$

where

$$L'_T = |\vec{L}'_T|, \quad P'_T = |\vec{P}'_T|, \quad y_L = \frac{1}{2} \log \left( \frac{L'^+}{L'^-} \right), \quad y_P = \frac{1}{2} \log \left( \frac{P'^+}{P'^-} \right)$$



## Virtual Photon-Hadron Frame

- Traditionally Breit frame is photon-hadron frame
- Lepton radiation makes frame determination ambiguous
- All historical factorization formula defined in photon hadrone frame
- Introduce **virtual** photon-hadron frame which is determined by a given pair of  $\xi, \zeta$  under one-photon exchange approximation

$$\begin{aligned}x_B &= \frac{Q^2}{2P \cdot q} &\Rightarrow & \hat{x}_B = \frac{\hat{Q}^2}{2P \cdot \hat{q}} \\y &= \frac{P \cdot q}{P \cdot L} &\Rightarrow & \hat{y} = \frac{P \cdot \hat{q}}{P \cdot L} \\z_h &= \frac{P \cdot P'}{P \cdot q} &\Rightarrow & \hat{z}_h = \frac{P \cdot P'}{P \cdot \hat{q}}\end{aligned}$$



# Transformation Between Frames

- Given generic vector in Breit frame

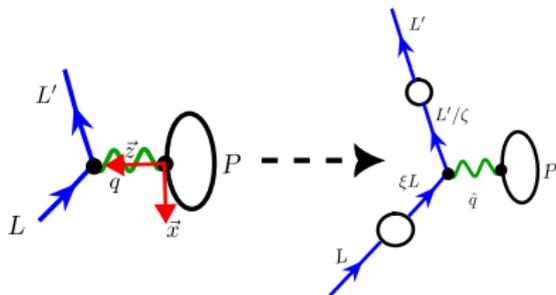
$$x^\mu = (x^+, x^-, \vec{x}_\perp = (x^1, x^2))$$

- Lorentz boost ( $\phi$ ) to hadron rest frame then rotate ( $\theta(\xi, \zeta)$ ) around  $y$ -axis ( $\perp$  to lepton plane) so  $\gamma$  along  $z$ -axis

$$\tilde{x}^\sigma = R_\nu^\sigma \Lambda_\mu^\nu x^\mu$$

$$y \Rightarrow \tilde{y} = y$$

$$\begin{pmatrix} x \\ z \end{pmatrix} \Rightarrow \begin{pmatrix} \tilde{x} \\ \tilde{z} \end{pmatrix} = R(\theta(\xi, \zeta)) \begin{pmatrix} x \\ z \end{pmatrix}$$



# Helicity Basis

- Vectors defined similar to (Ji et al., 2006)
- Work with helicity-based hadronic structure functions (similar to lepton case in (Liu et al., 2021))

$$\hat{W}^{\mu\nu} = \frac{1}{2} \underbrace{(\hat{X}^\mu \hat{X}^\nu + \hat{Y}^\mu \hat{Y}^\nu)}_{\text{Transverse}} \hat{H}_{TT} + \underbrace{(\hat{T}^\mu \hat{T}^\nu \hat{H}_L + (\hat{T}^\mu \hat{X}^\nu + \hat{X}^\mu \hat{T}^\nu) \hat{H}_\Delta + \dots)}_{\text{other, } \mathcal{O}(\alpha_S^2)}$$

$$\hat{Z}^\mu = -\frac{\hat{q}^\mu}{\hat{Q}}$$

$$\hat{T}^\mu = \left(\frac{1}{\hat{Q}}\right) (\hat{q}^\mu + 2\hat{x}_B P^\mu)$$

$$\hat{X}^\mu = \left(\frac{1}{\vec{\hat{q}}_T}\right) \left(\frac{(P')^\mu}{\hat{z}_h} - \hat{q}^\mu - \left(1 + \frac{\vec{\hat{q}}_T^2}{\hat{Q}^2}\right) \hat{x}_B P^\mu\right)$$

$$\hat{Y}^\mu = \epsilon^{\mu\nu\alpha\beta} \hat{Z}_\nu \hat{T}_\alpha \hat{X}_\beta$$



# Structure Functions

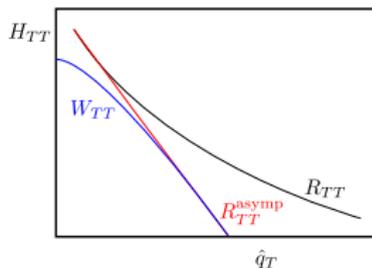
- For the fixed order calculation, where  $W_{\mu\nu}$  projects the hadron to the partonic states,

$$H_{TT} = \frac{1}{2}(X^\mu X^\nu + Y^\mu Y^\nu)W_{\mu\nu}$$

- where the real and virtual parts as power series of  $\alpha_S$  form  $W_{\mu\nu} = R_{\mu\nu} + V_{\mu\nu}$
- These also depend on  $\vec{q}_T$ :
  - When  $\vec{q}_T \sim \hat{Q}$ , fixed order is under control (no large logarithms)
  - when  $\vec{q}_T \gg \hat{Q}$  or  $\vec{q}_T \ll \hat{Q}$ , large logarithm divergences (focus on latter)

$$H_{TT}(\vec{q}_T) = W_{TT}(\text{small } \vec{q}_T) + Y_{TT}(\text{large } \vec{q}_T)$$

$$Y_{TT} = R_{TT} - R_{TT}^{\text{asympt}}$$



## CSS-like Formalism

- Better to work in Fourier Transform of  $\vec{q}_T$  to  $\vec{b}_T$  space (as in the CSS formalism (Collins et al., 1985))

$$\tilde{W}_{TT}(\vec{b}_T, Q) = \int d^2\vec{q}_T e^{i\vec{q}_T \cdot \vec{b}_T} W_{TT}(\vec{q}_T, Q) = e^{-S} [C_f \otimes f] \otimes [C_D \otimes D]$$

- Expand resummed expression for  $\tilde{W}_{TT}(\vec{b}_T)$  up  $\mathcal{O}(\alpha_S)$

$$\tilde{W}_{TT} = C_f^{(1)} C_D^{(0)} S^{(0)} + C_f^{(0)} C_D^{(1)} S^{(0)} + C_f^{(0)} C_D^{(0)} S^{(1)}$$

where

$$S = 1 - \frac{\alpha}{\pi} \left[ \frac{1}{2} A^{(1)} \ln^2 \frac{\nu_Q^2}{\mu_b^2} + B^{(1)} \ln \frac{\nu_Q^2}{\mu_b^2} \right]$$

- Important for defining hadronic structure function, which are related to PDFs



## NLO Perturbative Coefficients

$$\begin{aligned} \tilde{W}_{TT} = & \frac{4e_Q^2 \alpha_S}{3(2\pi)} \left[ \left( \left( \frac{2\lambda^2 - 2\lambda + 1}{\lambda} \right) \frac{\delta(1-\eta)}{(1-\lambda)_+} + \left( \frac{2\eta^2 - 2\eta + 1}{\eta} \right) \frac{\delta(1-\lambda)}{(1-\eta)_+} \right) \right. \\ & + \left( \left( \frac{\lambda^2 + 1}{\lambda} \right) \frac{\delta(1-\eta)}{(1-\lambda)_+} + \left( \frac{\eta^2 + 1}{\eta} \right) \frac{\delta(1-\lambda)}{(1-\eta)_+} \right) \left[ -\frac{1}{\epsilon} - \ln \frac{\mu_{\overline{MS}}^2}{\mu_b^2} \right] \\ & \left. + 2\delta(1-\lambda)\delta(1-\eta) \left( -\frac{2}{\epsilon} + \frac{1}{2} \left( -\ln^2 \frac{\mu_b^2}{\nu_Q^2} - 3 \ln \frac{\mu_b^2}{\nu_Q^2} - 3 \ln \frac{\mu_{\overline{MS}}^2}{\mu_b^2} - 4 \right) \right) \right] \quad (3) \end{aligned}$$

$$A^{(1)} = 1 \quad (4)$$

$$B^{(1)} = -\frac{3}{2} \quad (5)$$

$$C_f^{(1)}(\lambda) = \frac{1}{2\lambda}(1-2\lambda) - \frac{1}{\lambda} \left( \frac{\lambda^2 + 1}{1-\lambda} \right)_+ \ln \frac{\mu_{\overline{MS}}}{\mu_b} - \delta(1-\lambda) \quad (6)$$

$$C_D^{(1)}(\eta) = \frac{1}{2\eta}(1-2\eta) - \frac{1}{\eta} \left( \frac{\eta^2 + 1}{1-\eta} \right)_+ \ln \frac{\mu_{\overline{MS}}}{\mu_b} - \delta(1-\eta) \quad (7)$$



## Conclusions

- Applied joint factorization scheme for SIDIS process
- Derived general Lorentz transformation between virtual photon-hadron frame and lepton-hadron lab frame (or Lab-Breit frame)
- Derived perturbative coefficients in the CSS formalism in the virtual photon-hadron frame
- Next Steps:
  - Prepare code to make predictions for unpolarized and polarized PDFs from SIDIS in CSS approach
  - Expand to study the TMD approach
  - Use known PDFs to test sensitivities of  $\Delta G$  on  $q$  and  $G$  at JLab 12 and EIC



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